Albert-Ludwigs-Universität Freiburg

Particle Physics II

Markus Schumacher, Anna Kopp, Stan Lai

Problem Set IV

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In-class exercises

Exercise 23 Completeness of Dirac Equation Solutions

Given these four solutions to the Dirac Equation:

$$u^{(1)} = N \begin{pmatrix} 1\\ 0\\ \frac{p_z}{E+m}\\ \frac{p_x+ip_y}{E+m} \end{pmatrix}, u^{(2)} = N \begin{pmatrix} 0\\ 1\\ \frac{p_x-ip_y}{E+m}\\ \frac{-p_z}{E+m} \end{pmatrix}, v^{(1)} = N \begin{pmatrix} \frac{p_x-ip_y}{E+m}\\ -\frac{-p_z}{E+m}\\ 0\\ 1 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m}\\ \frac{p_x+ip_y}{E+m}\\ 1\\ 0 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m}\\ -\frac{p_z}{E+m}\\ 1\\ 0 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m}\\ -\frac{p_z}{E+m}\\ 1\\ 0 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m}\\ -\frac{p_z}{E+m}\\ 1\\ 0 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m}\\ -\frac{p_z}{E+m}\\ 1\\ 0 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m}\\ -\frac{p_z}{E+m}\\ 1\\ 0 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m}\\ -\frac{p_z}{E+m}\\ 1\\ 0 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m}\\ -\frac{p_z}{E+m}\\ -\frac{p_z}{E+m}\\ 1 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m}\\ -\frac{p_z}{E+m}\\ -\frac{p_z}{E+m$$

with $N = \sqrt{E+m}$, show that the completeness relation holds for these spinors, namely that:

$$\sum_{s=1,2} u^{(s)} \bar{u}^{(s)} = (\gamma^{\mu} p_{\mu} + m), \sum_{s=1,2} v^{(s)} \bar{v}^{(s)} = (\gamma^{\mu} p_{\mu} - m)$$

Exercise 24 Møller Scattering

Bhabha scattering is name for the QED process $e^+e^- \rightarrow e^+e^-$, while Møller scattering is the name for the process $e^-e^- \rightarrow e^-e^-$.

- (a) Draw the Feynman diagrams for Møller scattering.
- (b) What is the relative sign between the two diagrams? (Important for the interference term in the Matrix element squared)
- (c) In the lectures, it was shown that the matrix element squared for Bhabha scattering written in terms of Mandelstam variables, can be simplified to:
 (after evencing over initial state gring, and summing over the final state gring)

(after averaging over initial state spins, and summing over the final state spins)

$$|\bar{\mathcal{M}}|^2 = 2e^4 \left[\frac{u^2 + t^2}{s^2} + \frac{2u^2}{st} + \frac{u^2 + s^2}{t^2} \right]$$

Using crossing symmetry, arrive at the matrix element squared $|\overline{\mathcal{M}}|^2$ for Møller scattering. (also for averaging over initial state spins and summer over final state spins)

Homework

Exercise 25 More Møller Scattering

- (a) From the result in Problem 24c, write down the differential cross-section for Møller scattering. You can assume that the electrons are massless (i.e. $\sqrt{s} \gg m_e$). Don't forget that you have identical particles in the initial and final states!
- (b) Show that if we substitute $\theta \to \pi \theta$, that the result remains unchanged. Why is that?

Exercise 26 The Rutherford Limit for Mott Scattering

Mott scattering entails the process $e^-X \to e^-X$, where X has a mass much larger than the electron energy (could be a muon or a proton). From the lectures, we have calculated the matrix element squared (averaging over initial state spins, and summing over final state spins):

$$|\bar{\mathcal{M}}_{Mott}|^2 = \frac{8e^4}{(p_1 - p_3)^4} \left[(p_1 \cdot p_2)(p_3 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3) - m_X^2(p_1 \cdot p_3) - m_e^2(p_2 \cdot p_4) + 2m_e^2 m_X^2 \right].$$

(a) In the rest-frame of the heavy particle X, the heavy particle is stationary before and after the collision $(m_X >> |\vec{p_e}|)$. Show the the Mott scattering scross-section in this frame-of-reference is

$$\left(\frac{d\sigma}{d\Omega}\right)_{Mott} = \frac{\alpha^2}{4|\vec{p_e}|^4 \sin^4(\theta/2)} \left[m_e^2 + \vec{p}^2 \cos^2(\theta/2)\right].$$

(b) Consider the case where the incident electron is non-relativistic $|\vec{p}_e|^2 \ll m_e^2$. Show that the differential cross section reduces to the Rutherford scattering formula:

$$\left(\frac{d\sigma}{d\Omega}\right)_{Rutherford} = \frac{\alpha^2}{4m_e^2 v^4 \sin^4(\theta/2)}$$

Exercise 27 Threshold Behaviour for $e^+e^- \rightarrow \mu^+\mu^-$

The unpolarized matrix element for $e^+e^- \rightarrow \mu^+\mu^-$ or for any fermion-antifermion final state with a mass much larger than the electron mass can be expressed as:

$$|\bar{\mathcal{M}}|^2 = \frac{2e^4}{(p_1 \cdot p_2)^2} \left[(p_1 \cdot p_3)(p_2 \cdot p_4) + (p_1 \cdot p_4)(p_2 \cdot p_3) + m_{\mu}^2(p_1 \cdot p_2) \right].$$

Now we want to find the behaviour of the cross-sections for $e^+e^- \rightarrow \mu^+\mu^-$ at threshold, so we should not neglect the muon mass in the following:

- (a) Express the differential cross-section in terms of the Mandelstam variable $s = E_{CM}^2$ and the relativistic velocity $\beta = \frac{|\vec{p}|}{E}$.
- (b) Integrate this to obtain an expression of the total cross-section.

6 Points

7 Points

7 Points