

# Triple Gauge Couplings and Quartic Gauge Couplings

## Particle Physics at the LHC

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# QED Lagrangian

Quantum electrodynamics described by

$$\mathcal{L} = \bar{\psi}(i\gamma^\mu \partial_\mu - m)\psi - Q\bar{\psi}\gamma^\mu\psi A_\mu - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}.$$

free fermion propagation ( $\mathcal{L}_0$ )

photon kinetic energy

fermion photon interaction

Theory invariant under local phase transformation ( $U(1)$ )

$$\psi \rightarrow \psi' = e^{iQ\alpha(x)}\psi.$$

$$(A_\mu \rightarrow A'_\mu = A_\mu + \partial_\mu\chi)$$

# Electroweak Lagrangian

Theory invariant under  $SU(2)_L \otimes U(1)$

$$\psi \rightarrow \psi' = e^{\frac{i}{2} \vec{\tau} \vec{\alpha}(x)} e^{i \frac{Y}{2} \beta(x)} \psi.$$

$$\mathcal{L} = \bar{\chi}_L \gamma^\mu \left( i \partial_\mu - g \frac{\vec{\tau}}{2} \vec{W}_\mu - g' \frac{Y}{2} B_\mu \right) \chi_L + \bar{\psi}_R \left( i \partial_\mu - g' \frac{Y}{2} B_\mu \right) \psi_R + \mathcal{L}_{kin}$$

with gauge invariant term describing gauge field kinetics:

$$\mathcal{L}_{kin} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} \vec{W}_{\mu\nu} \vec{W}^{\mu\nu}$$

*resulting from non-abelian gauge structure*

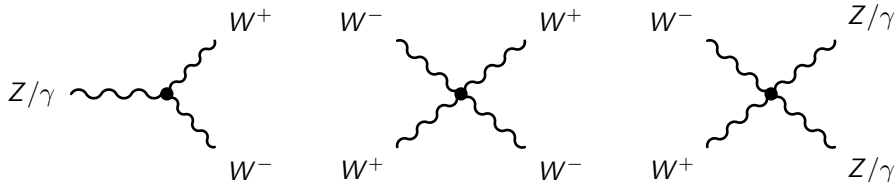
$$B_{\mu\nu} := \partial_\mu B_\nu - \partial_\nu B_\mu \quad \vec{W}_{\mu\nu} := \partial_\mu \vec{W}_\nu - \partial_\nu \vec{W}_\mu + \overbrace{g \vec{W}_\mu \times \vec{W}_\nu}$$

# Gauge boson self-coupling in the electroweak theory

Cubic and quartic interaction terms resulting from  $\mathcal{L}_{kin}$ :

$$\mathcal{L}_3 = -ie \cot \theta_W \left[ (\partial^\mu W^\nu - \partial^\nu W^\mu) W_\mu^\dagger Z_\nu - (\partial^\mu W^{\nu\dagger} - \partial^\nu W^{\mu\dagger}) W_\mu Z_\nu + W_\mu W_\nu^\dagger (\partial^\mu Z^\nu - \partial^\nu Z^\mu) \right] \\ - ie \left[ (\partial^\mu W^\nu - \partial^\nu W^\mu) W_\mu^\dagger A_\nu - (\partial^\mu W^{\nu\dagger} - \partial^\nu W^{\mu\dagger}) W_\mu A_\nu + W_\mu W_\nu^\dagger (\partial^\mu A^\nu - \partial^\nu A^\mu) \right]$$

$$\mathcal{L}_4 = -\frac{e^2}{2 \sin^2 \theta_W} \left[ (W_\mu^\dagger W^\mu)^2 - W_\mu^\dagger W^{\mu\dagger} W_\nu W^\nu \right] - e^2 \cot^2 \theta_W \left[ W_\mu^\dagger W^\mu Z_\nu Z^\nu - W_\mu^\dagger Z^\mu W_\nu Z^\nu \right] \\ - e \cot \theta_W \left[ 2W_\mu^\dagger W^\mu Z_\nu A^\nu - W_\mu^\dagger Z^\mu W_\nu A^\nu - W_\mu^\dagger A^\mu W_\nu Z^\nu \right] \\ - e^2 \left[ W_\mu^\dagger W^\mu A_\nu A^\nu - W_\mu^\dagger A^\mu W_\nu A^\nu \right]$$



# Effective field theory (EFT)

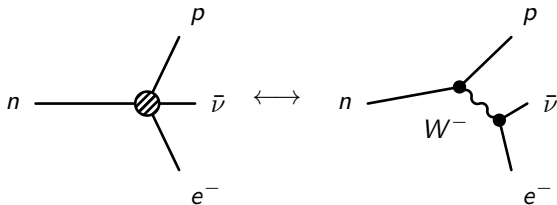
**Assumption:** New physics separated by different (energy) scale  $\Lambda$  from accessible region. ( $s \ll \Lambda^2$ )

→ Describe observations by parametrized, most general Lagrangian which recovers the Standard Model in the limit  $\Lambda \rightarrow \infty$ .

(Unitarity should be preserved.)

→ Model independent approach for physics BSM.

**Example:** Fermi theory of  $\beta$  decay (quartic coupling with coupling strength  $G_F$ , describing two weak interaction vertices at low energies  $s \ll M_W$ .)



# Limits on anomalous neutral triple gauge couplings

# Neutral triple gauge couplings (nTGC)

Forbidden in Standard Model but possibly realized as “anomalous” coupling (in EFT approach).

Construct Lagrangian from general process properties:

- 9 total helicity states but only 7 valid (angular momentum conservation)
- bose statistics have to be respected

→ Effective ZZV Lagrangian:

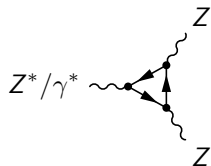
$$\mathcal{L}_{ZZV} = \frac{e^2}{M_Z^2} \left( \begin{array}{l} \text{violating CP invariance} \\ - [\mathbf{f}_4^\gamma (\partial_\mu F^{\mu\nu}) + \mathbf{f}_4^Z (\partial_\mu Z^{\mu\nu})] Z_\sigma (\partial^\sigma Z_\nu) \\ \\ \text{conserving CP invariance} \rightarrow - [\mathbf{f}_5^\gamma (\partial^\rho F_{\rho\lambda}) + \mathbf{f}_5^Z (\partial^\rho Z_{\rho\lambda})] \tilde{Z}^{\lambda\xi} Z_\xi \end{array} \right)$$



# SM and some new physics contributions to nTGC

## SM

NLO and higher order contributions  
(only CP conserving  $f_5^V$ )



## MSSM

1-loop (and higher order) contributions from charginos and neutralinos

## New bosons

CP violating coupling  $f_4^V$  sensitive to two-Higgs-doublet model in 1-loop corrections.

# Limits on anomalous nTGC from $ZZ$ production in pp collisions

## Strategy:

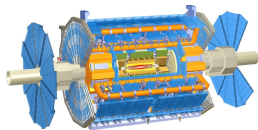
Obtain limits on anomalous  $ZZZ$  and  $ZZ\gamma$  couplings from differential

cross section  $\frac{d\sigma_{ZZ}}{dp_T^Z}$ .

Analyse  $ZZ$  signal channels:  $ZZ^{(*)} \rightarrow l^+l^-l^+l^-$  and  $ZZ \rightarrow l^+l^-\nu\bar{\nu}$

# Extended-lepton selection

**Aim:** Increase selection acceptance in the  $ZZ^{(*)} \rightarrow l^+l^-l^+l^-$  channel by using leptons which are normally not used due to detector geometry.



## Forward spectrometer muons

Muons outside the nominal ID range with  $2.5 < |\eta| < 2.7$ .

Are required to have a) full track in muon spectrometer, b)  $p_T > 10$  GeV and c)

$\sum E_T$  of calorimeter deposits inside  $\Delta R = 0.2$  smaller than 15% of muon  $p_T$ .

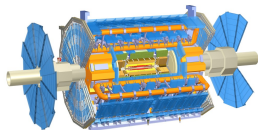
## Calorimeter-tagged muons

Muons in the muon spectrometer limited coverage range  $|\eta| < 0.1$ .

Are required to a) have calorimeter deposit consistent with muon which is matched to ID track b) have  $p_T > 20$  GeV and c) fulfill same impact parameter and isolation criteria as “standard muons”.

# Extended-lepton selection

**Aim:** Increase selection acceptance in the  $ZZ^{(*)} \rightarrow l^+l^-l^+l^-$  channel by using leptons which are normally not used due to detector geometry.



## Calorimeter-only electrons

Electrons outside the ID range with  $2.5 < |\eta| < 3.16$ .

Are required to a) have  $p_T > 20$  GeV and b) pass the tight identification requirement.

$p_T$  is calculated from calorimeter energy and electron direction.

Charge is assigned depending on the charge of the other electron(s).

At most one lepton from each extended category!

# Signal region definitions

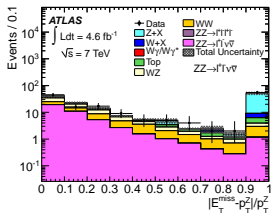
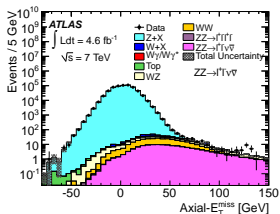
$ZZ^{(*)} \rightarrow l^+l^-l^+l^-$ :

- exactly 4 isolated leptons, two same-flavour opposite charge pairs ( $e^+e^-e^+e^-$ ,  $e^+e^-\mu^+\mu^-$  or  $\mu^+\mu^-\mu^+\mu^-$ )
- $\Delta R(l_1, l_2) > 0.2$
- ambiguity in lepton combinations removed by choosing combination with lowest  $|m_{l+l-} - M_Z|$
- at least one lepton pair fulfills  $66 < m_{l+l-} < 116$  GeV (the other  $m_{l+l-} > 20$  GeV)

## Signal region definitions

$$ZZ \rightarrow l^+l^-\nu\bar{\nu}:$$

- exactly 2 leptons of same flavour with  $p_T > 20$  GeV
- $\Delta R(l_1, l_2) > 0.3$
- $76 < m_{l^+l^-} < 106$  GeV
- axial- $E_T^{miss} = -\vec{E}_T^{miss} \cdot \vec{p}_T^Z / p_T^Z > 75$  GeV and  $|E_T^{miss} - p_T^Z| / p_T^Z < 0.4$
- jet veto
- no additional lepton with  $10 < p_T \leq 20$  GeV



Background estimation for the  $ZZ^{(*)} \rightarrow l^+l^-l^+l^-$  channel

Background estimated via data-driven (dd.) method

$$N(BG) = [N(IIIj) - N(ZZ)] \times f - N(IIjj) \times f^2,$$

where

- $N(IIIj)$  = number of events with 3 leptons and 1 lepton-like jet satisfying all selection criteria
- $N(IIjj)$  = number of events with 2 leptons and 2 lepton-like jet satisfying all selection criteria
- $N(ZZ)$  = MC estimate for real leptons classified as lepton-like jet
- $f$  = ratio of the probability for a non-lepton to satisfy the full lepton selection criteria to the probability for a non-lepton to satisfy the lepton-like jet criteria.

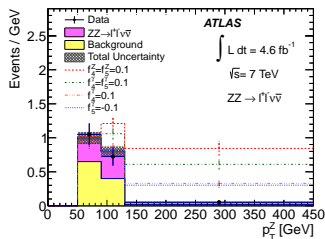
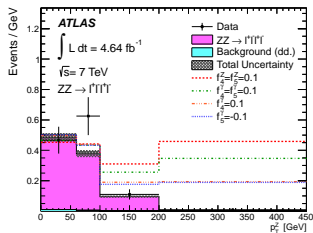
# Determining limits on anomalous nTGC

- Couplings are parametrized in form-factor approach

$$f_i^V = \frac{1}{(1 + \hat{s}/\Lambda^2)^n} f_{i,0}^V \xrightarrow{s \rightarrow \infty} 0$$

with  $n = 3$  and  $\Lambda = 3 \text{ TeV}$  to ensure unitarity is not violated at LHC energies.

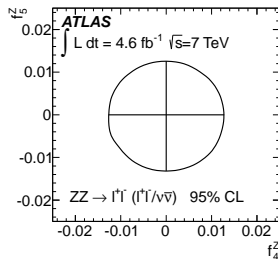
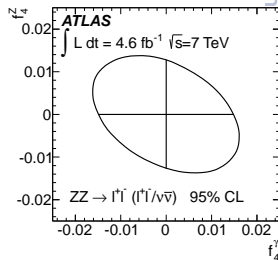
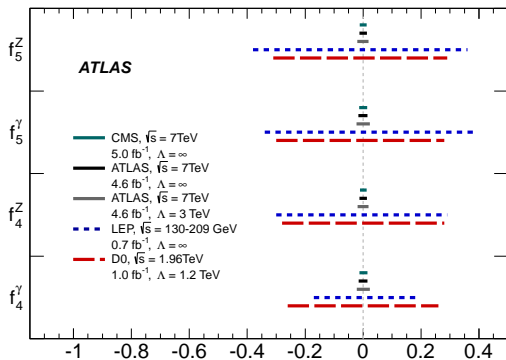
- To obtain simulated  $p_T^Z$  distributions for different  $f_i^V$  a reweighting method ( $|\mathcal{M}|^2/|\mathcal{M}_{SM}|^2$ ) is used.
- Dependency of couplings on the expected number of events in each  $p_T^Z$  bin is parametrized.
- Limits on couplings are obtained by using a maximum likelihood fit.





# Limits on anomalous nTGC

The limit(s) on the coupling(s) is/are obtained assuming all other couplings are zero (as in SM).



# Limits on charged triple gauge couplings

# Charged triple gauge couplings (cTGC)

Already realized as  $ZWW$  and  $\gamma WW$  in Standard Model but further coupling contributions from new physics possible in EFT approach.

Construct Lagrangian from general process properties:

- 9 total helicity states but only 7 valid (angular momentum conservation)
- demand C and P conservation

→ Effective WWV Lagrangian ( $V = Z$  and  $\gamma$ ):

$$\mathcal{L}_{WWV} = ig_{WWW} \left( g_1^V (W_{\mu\nu}^+ W^{-\mu} - W^{+\mu} W_{\mu\nu}^-) V^\nu + \kappa_V W_\mu^+ W_\nu^- V^{\mu\nu} + \frac{\lambda_V}{M_W^2} V^{\mu\nu} W_\nu^{+\rho} W_{\rho\mu}^- \right)$$

with  $g_1^\gamma = 1$ .

# The values of cTGC

In the SM the general cTGC couplings are given by

$$g_1^Z = \kappa_Z = \kappa_\gamma = 1$$
$$\lambda_Z = \lambda_\gamma = 0.$$

Often the differences from the SM

$$\Delta g_1^Z = g_1^Z - 1$$
$$\Delta \kappa_Z = \kappa_Z - 1$$
$$\Delta \kappa_\gamma = \kappa_\gamma - 1$$

and not the absolute values are denoted.

# Limits on anomalous charged TGC from $W^+W^-$ production in $pp$ collisions

## Analysis strategy:

- Select  $W^+W^-$  events by " $ll' + E_T^{miss}$ ".
- Measure differential cross section  $\frac{d\sigma_{WW}}{dp_T}$  in selection phase space region.
- Extract limits on couplings from differential cross section.

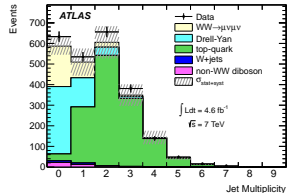
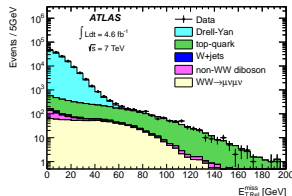
# Selection criteria

- two opposite charged leptons  
(at least one matched to trigger reconstructed lepton)  
( $\rightarrow$  3 channels:  $ee$ ,  $e\mu$ ,  $\mu\mu$ )
- cut on invariant lepton mass  $m_{ll'}$  and  $E_{T,Rel}^{miss}$ ,  
where

$$E_{T,Rel}^{miss} = \begin{cases} E_T^{miss} \times \sin(\Delta\phi) & \text{if } \Delta\phi^a < \pi/2 \\ E_T^{miss} & \text{if } \Delta\phi \geq \pi/2 \end{cases}$$

(to remove Drell-Yan background)

- jet veto
- $p_T(l'l') > 30$  GeV



<sup>a</sup> $\Delta\phi$  = azimuthal angle difference between  $E_T^{miss}$  and nearest lepton or jet

# “Coupling-scenarios” investigated

## equal coupling scenario

$$\Delta\kappa_Z = \Delta\kappa_\gamma, \lambda_Z = \lambda_\gamma \text{ and } g_1^Z = 1$$

## LEP scenario

$$\Delta\kappa_\gamma = \frac{\cos^2 \theta_W}{\sin^2 \theta_W} \left( \Delta g_1^Z - \Delta\kappa_Z \right) \text{ and } \lambda_Z = \lambda_\gamma$$

## HISZ scenario

$$\Delta g_1^Z = \frac{1}{\cos^2 \theta_W - \sin^2 \theta_W} \Delta\kappa_Z,$$

$$\Delta\kappa_\gamma = 2\Delta\kappa_Z \frac{\cos^2 \theta_W}{\cos^2 \theta_W - \sin^2 \theta_W} \text{ and } \lambda_Z = \lambda_\gamma$$

# Obtaining limits on cTGC

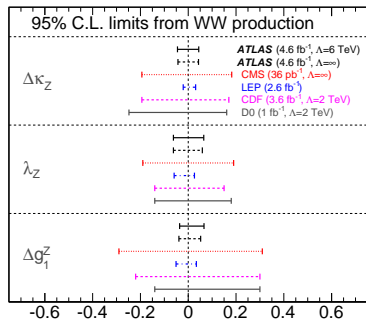
Similar to the  $pp \rightarrow ZZ$  analysis (form factor approach, reweighting method, dependency of  $p_T(l_1)$  on couplings).

Limits are obtained by the maximum likelihood principle.  
(Point in parameter space is discarded if the negative log-likelihood function increases by more than 1.92 units above the minimum.)



## Limits on cTGC

Scenario	Parameter	Expected	Observed	Expected	Observed
		( $\Lambda = 6$ TeV)	( $\Lambda = 6$ TeV)	( $\Lambda = \infty$ )	( $\Lambda = \infty$ )
LEP	$\Delta\kappa_Z$	[-0.043, 0.040]	[-0.045, 0.044]	[-0.039, 0.039]	[-0.043, 0.043]
	$\lambda_Z = \lambda_\gamma$	[-0.060, 0.062]	[-0.062, 0.065]	[-0.060, 0.056]	[-0.062, 0.059]
	$\Delta g_1^Z$	[-0.034, 0.062]	[-0.036, 0.066]	[-0.038, 0.047]	[-0.039, 0.052]
HISZ	$\Delta\kappa_Z$	[-0.040, 0.054]	[-0.039, 0.057]	[-0.037, 0.054]	[-0.036, 0.057]
	$\lambda_Z = \lambda_\gamma$	[-0.064, 0.062]	[-0.066, 0.065]	[-0.061, 0.060]	[-0.063, 0.063]
Equal Couplings	$\Delta\kappa_Z$	[-0.058, 0.089]	[-0.061, 0.093]	[-0.057, 0.080]	[-0.061, 0.083]
	$\lambda_Z = \lambda_\gamma$	[-0.060, 0.062]	[-0.062, 0.065]	[-0.060, 0.056]	[-0.062, 0.059]



Parameter	Expected	Observed
	( $\Lambda = \infty$ )	( $\Lambda = \infty$ )
$\Delta\kappa_Z$	[-0.077, 0.086]	[-0.078, 0.092]
$\lambda_Z$	[-0.071, 0.069]	[-0.074, 0.073]
$\lambda_\gamma$	[-0.144, 0.135]	[-0.152, 0.146]
$\Delta g_1^Z$	[-0.449, 0.546]	[-0.373, 0.562]
$\Delta\kappa_\gamma$	[-0.128, 0.176]	[-0.135, 0.190]

# Limits on anomalous quartic gauge couplings

# Quartic gauge couplings (QGC)

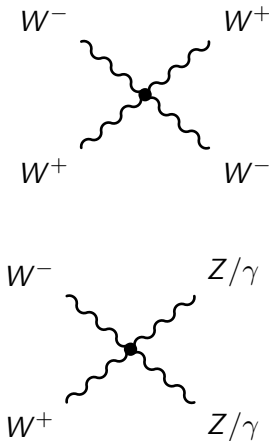
In SM realized as  $WWWW$ ,  $WWZZ$ ,  $WW\gamma Z$  and  $WW\gamma\gamma$ . Further contributions from physics BSM in EFT approach possible.

Construct Lagrangian from following constraints:

- Consider couplings to longitudinal polarization degree of gauge bosons only.
- Custodial symmetry

$$\rho = \left( \frac{M_W}{M_Z \cos \theta_W} \right)^2 = 1$$

should be respected.



# Quartic gauge couplings (QGC)

→ Two effective anomalous vector boson scattering (VBS)  
Lagrangian terms:

$$\mathcal{L}_4 = \frac{\alpha_4}{16\pi^2} \text{Tr}(V_\mu V_\nu) \text{Tr}(V^\mu V^\nu)$$

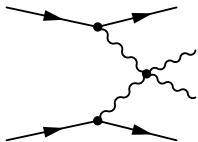
$$\mathcal{L}_5 = \frac{\alpha_5}{16\pi^2} \text{Tr}(V_\mu V^\mu) \text{Tr}(V_\nu V^\nu)$$

with  $V_\mu = -igW_\nu + ig'B_\mu$  (in unitary gauge)

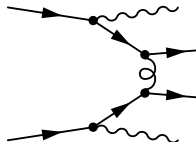
Limits on aQGC from  $W^\pm W^\pm jj$  production in  $pp$  collisions

“Evidence of the electroweak production of  $W^\pm W^\pm jj$  in  $pp$  collisions at  $\sqrt{s} = 8$  TeV with the ATLAS detector” (*ATLAS-CONF-2014-013*)

- Published in March 2014.
- *First analysis* being able to give direct constraints on QGC from VBS.
- Limits on aQGC determined from measured cross section in signal region phase space.



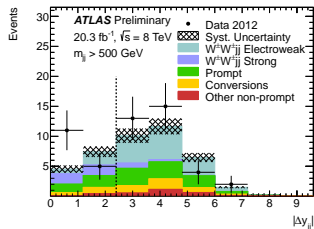
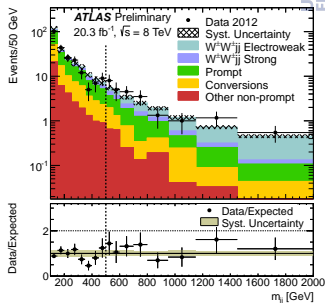
$W^\pm W^\pm jj$  electroweak



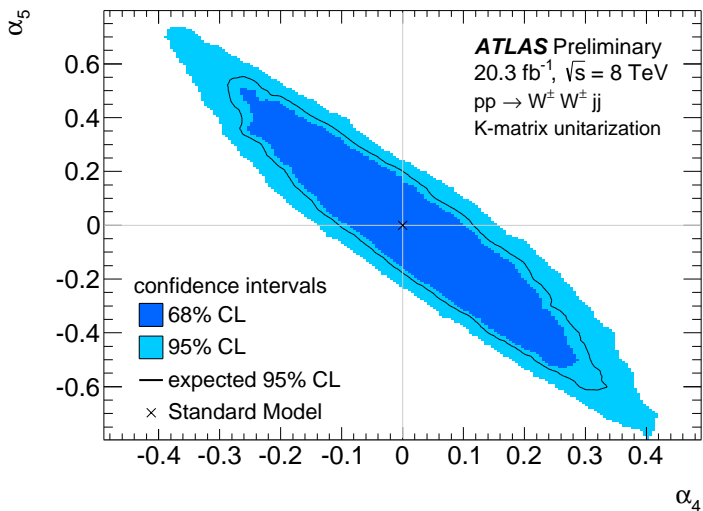
$W^\pm W^\pm jj$  strong

## Signal region definition

- 2 same electric charge leptons
- 2 jets with  $p_T > 30$  GeV &  $|\eta| < 4.5$
- additional loose lepton veto
- $m_{ll} > 20$  GeV
- $|m_{ee} - M_Z| > 10$  GeV
- $E_T^{miss} > 40$  GeV
- b-jet veto
- $m_{jj} > 500$  GeV
- $|\Delta y_{jj}| > 2.4$



## Limits on aQGC



# Summary & Outlook

## Summary

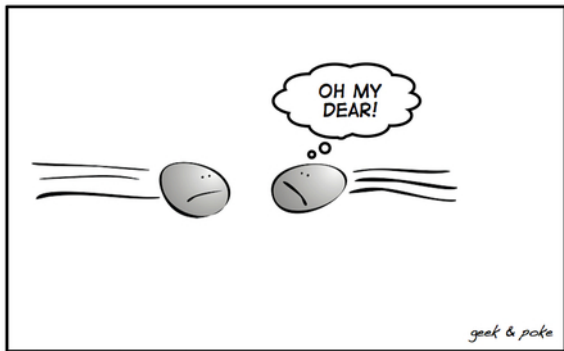
- Effective field theory approach yields possibility to probe physics BSM in model independent way.
- LHC experiments could improve the limits on anomalous nTGC from LEP and reach a similar precision on limits on cTGC.
- New analysis (03/2014) gives first direct constraints on aQGC from  $pp \rightarrow W^\pm W^\pm jj$ .
- No physics BSM observed! (But higher precision may show contributions!)

## Outlook

- Run 2 of LHC will probably increase sensitivity.
- Triboson production analysis (yielding limits on aQGC) may be realizable for new phase.



# The END



*LATELY INSIDE THE LHC:  
2 PROTONS 0.000000000000000001 SEC BEFORE THE COLLISION*