

# Particle Physics II

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## Problem Set II

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### In-class exercises

#### Exercise 8 Conservation of Total Angular Momentum

- (a) The Hamiltonian in relativistic quantum mechanics is just  $H \equiv p^0$ . Write the Hamiltonian  $H$  in terms of 3-momentum and mass operators, by rearranging the Dirac equation.  
(Remember the quantum mechanical momentum operator for momentum is given by  $p_\mu = i\partial_\mu$ .)
- (b) Show that the Hamiltonian ( $H \equiv p^0$ ) for the Dirac Equation does not commute with the orbital angular momentum operator:

$$[H, \vec{L}] = -i\gamma^0(\vec{\gamma} \times \vec{p}) \neq 0.$$

(Since  $\vec{L} = \vec{r} \times \vec{p}$ , we can write  $L^i = \epsilon^{ijk} r_j p_k$ .)

- (c) Show also that this Hamiltonian does not commute with the spin operator:

$$[H, \vec{S}] = i\gamma^0(\vec{\gamma} \times \vec{p}).$$

(Recall that  $\vec{S} = \vec{\Sigma}/2$  whereby  $\vec{\Sigma}$  are just the Pauli spin matrices  $\sigma_x, \sigma_y, \sigma_z$ .)

- (d) However, show that as a consequence, the total angular momentum  $\vec{J} = \vec{L} + \vec{S}$  is conserved.

*Hint:* recall that the commutation relation between position and momentum operators is:  $[x_k, p_l] = i\delta_{kl}$ .

#### Exercise 9 Completeness of Dirac Equation Solutions

Given these four solutions to the Dirac Equation:

$$u^{(1)} = N \begin{pmatrix} 1 \\ 0 \\ \frac{p_z}{E+m} \\ \frac{p_x + ip_y}{E+m} \end{pmatrix}, u^{(2)} = N \begin{pmatrix} 0 \\ 1 \\ \frac{p_x - ip_y}{E+m} \\ \frac{-p_z}{E+m} \end{pmatrix}, v^{(1)} = N \begin{pmatrix} \frac{p_x - ip_y}{E+m} \\ \frac{-p_z}{E+m} \\ 0 \\ 1 \end{pmatrix}, v^{(2)} = -N \begin{pmatrix} \frac{p_z}{E+m} \\ \frac{p_x + ip_y}{E+m} \\ 1 \\ 0 \end{pmatrix},$$

with  $N = \sqrt{E+m}$ , show that the completeness relation holds for these spinors, namely that:

$$\sum_{s=1,2} u^{(s)} \bar{u}^{(s)} = (\gamma^\mu p_\mu + m), \quad \sum_{s=1,2} v^{(s)} \bar{v}^{(s)} = (\gamma^\mu p_\mu - m).$$

#### Exercise 10 Adjunct Dirac Equation

Given the Dirac Equation

$$i\gamma^\mu \partial_\mu \psi - m\psi = 0,$$

show that the adjunct spinor  $\bar{\psi} \equiv \psi^\dagger \gamma^0$  satisfies the adjunct Dirac Equation

$$i\partial_\mu \bar{\psi} \gamma^\mu + m\bar{\psi} = 0.$$

**Exercise 11** *Target Scattering Energy Threshold*

If a particle  $A$ , with energy  $E$  hits a particle  $B$  at rest, and produces  $n$  particles  $C_1, C_2, \dots, C_n$  with masses  $m_1, m_2, \dots, m_n$ , what is the energy threshold (minimum incident energy  $E_{min}$ ) for this process to occur (in terms of  $m_A, m_B, m_1, \dots, m_n$ )?

## Homework

### Exercise 12 *Dirac and Klein-Gordon Equations*

4 Points

Show that a spinor  $\psi$  satisfying the Dirac equation  $(i\gamma^\mu\partial_\mu - m)\psi = 0$  also satisfies the Klein-Gordon equation  $(\partial^\mu\partial_\mu + m^2)\psi = 0$ .

*Hint:* Operate on the left for both sides of the Dirac equation with the operator  $\gamma^\nu\partial_\nu$ .

### Exercise 13 *Dirac Matrix Properties*

4 Points

Show the following are true, using the Dirac Representation of  $\gamma$ -Matrices:

(a)  $(\gamma^0)^2 = 1$

(b)  $(\gamma^k)^2 = -1$  (for  $k = 1, 2, 3$ )

(c)  $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}$

Show that (c) is also true for the Weyl representation.

What does  $\gamma^5$  look like in the Dirac representation and the Weyl representation?

### Exercise 14 *Orthogonality of Dirac Equation Solutions*

4 Points

For the above solutions to the Dirac equation, show that  $u^{(1)}$  and  $u^{(2)}$  are orthogonal, as are  $v^{(1)}$  and  $v^{(2)}$ . Show however, that  $u^{(1)}$  and  $v^{(1)}$  are not orthogonal.

*Hint:* Show for instance that  $u^{(1)\dagger}u^{(2)} = 0$ .

### Exercise 15 *Helicity in the Dirac Equation*

4 Points

The helicity operator is given by  $\hat{h} = \frac{\vec{\sigma}\cdot\vec{p}}{2|\vec{p}|}$ . Show that helicity is conserved in the Dirac equation.

*Hint:* Show that the helicity operator commutes with the Hamiltonian  $H \equiv p^0$ .

### Exercise 16 *Charge Conjugation and Time Reversal*

4 Points

The charge conjugation operator  $C$  transforms spinors via:  $\psi \rightarrow C(\psi) = \psi' = i\gamma^2\psi^*$ . What do the charge conjugates of the spinors  $v^{(1)}$  and  $v^{(2)}$  look like?

The time reversal operator  $T$  transforms spinors via:  $\psi \rightarrow T(\psi) = \psi' = i\gamma^1\gamma^3\psi^*$ . What does the time reversed spinor of  $u^{(1)}e^{-ip\cdot x}$  look like?